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Dr. Ryszard Gajewski
US Dept of Energy/AEP

by FAX

August 31, 1989

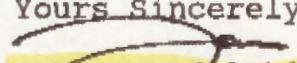
Dear Ryszard:

Enclosed please find a preliminary note on our understanding of the question of tritium production in Cold Fusion.

It would be very kind of you if it would be possible for you to forward a copy of this note to ER-6, 3F-061 (William Woodard) as per their request of 9 August. As I have no FAX number for Dr. Bigeleisen, I also ask them to attend to any further distribution of this preliminary document, intended only for INTERNAL use by the committee.

Please limit distribution to official use only until submission for publication.

Yours Sincerely


Johann Rafelski
Professor of Physics

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PRELIMINARY AZPH-TH/89-...

Comment on Tritium Production in Cold Fusion in Condensed Matter

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August 29, 1989

Following on the first report of the observation of cold fusion neutrons in laboratory, [1,2], Fleischmann *et al*[3] have also claimed that anomalous amounts of tritium may be produced in this way. While their other measurements of heat and million times higher neutron yield than that of Jones' have been repeatedly criticized, there have been recently several reports supporting the third claim that tritium is produced during the electrolysis of heavy water [4,5,6,7]. In particular, Packham *et al* have presented experimental data that places a lower limit on the tritium production rate in two such experiments with palladium cathode and nickel anode, vis. their runs labelled 'A2' and 'A7', for which we estimate a lower limit on the tritium production rate of $5.2 \times 10^6 \text{ ml}^{-1}$ of electrolyte s^{-1} and $2.1 \times 10^8 \text{ ml}^{-1}\text{s}^{-1}$ respectively. The same group concurrently measured the neutron flux during the experimental runs. In one electrolytical cell they obtain a neutron count of about one per second [8] which is about 10^6 less than the expected value derived from conventional deuteron-deuteron (dd) fusion, assuming symmetry between the neutron and tritium producing branches of the dd-fusion reaction. While the neutron count rate is consistent with the observation of 2.45 MeV fusion neutrons reported first by SE Jones *et al*[2], the reported tritium yield requires some phenomenon other than the one associated with the results of Jones *et al*. According to Wolf [8] it is very unlikely for the observed tritium abundance to have been from an impurity within the electrodes used in the experiments.

There are two points presented here in view of which make the tritium production perhaps even more mysterious than the excess heat reported. Firstly, the non-observation of neutrons severely limits the final-state triton energy due

to the possibility of a neutron producing deuterium-tritium (dt) secondary reaction. Secondly, the required smallness of this energy makes tritium production even less likely in an exit channel in which more than one charged particle is present (such as the proton-tritium (pt) exit channel of the dd reaction), due to the Coulomb suppression of such channels. These observations seem to imply that the only possible nuclear mechanism derives from some nuclear reaction other than dd, and we discuss all possible options without arriving at any satisfactory conclusion. Perhaps we should emphasize here that while we explore the consequences of energy and momentum conservation, which we presume to be sacred laws of physics, we are ignoring the fact that the lowest order quantum mechanical processes lead to rates for the processes we are discussing which are many orders of magnitude too small.

The constraint on the final state triton energy arises from the possibility of secondary fusion reactions being induced by high energy tritium produced in the 'cold fusion reaction'. It is well known that the dt reaction has an unusually large cross section that peaks at about 5 barn for a triton incident upon a deuterium target at an energy of about 70 KeV. The cross section is reduced to 2 barn at about 150 KeV, and falls off as $1/E_t$ at higher energies. If a high energy triton is produced in an environment containing substantial amounts of deuterium, it has a significant probability of undergoing a fusion reaction before being brought to rest. Specifically, the fraction of tritons of initial energy E that will fuse in a target is given by:

$$\frac{N_{fus}}{N_{inc}} \approx \int_0^E dE' \frac{\sigma_{fus}(E')}{S(E')} \quad (1)$$

$\sigma_{fus}(E)$ is the energy dependent fusion cross section. $S(E)$ is the stopping power of the target, and at low energies is generally dominated by electron ionization processes that dissipate the incident charged particle's momentum. We have computed this fraction in manner similar to our prior treatment of the suggested fractal fusion [9] as function of the incident triton energy, using experimental data for the fusion cross section for $d+t \rightarrow \alpha(2.8 \text{ MeV})+n(14.8 \text{ MeV})$ [10] and computing $S(E)$ for t incident upon a PdD target using a standardized approach [11]. Despite the relative smallness of the fusion fraction (7.9×10^{-12} fusions per triton at 10 KeV, 1.7×10^{-9} at 20 KeV, and increasing to 2.7×10^{-6} at 100 KeV) the absolute yield of secondary neutrons we should expect to see turns out to be forbiddingly high. For t-energies as small as 100 KeV we find that the secondary neutron production rate should be of the order of 14 neu-

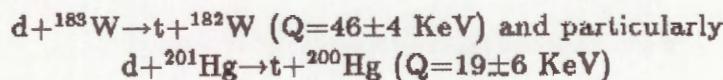
trons ml^{-1} electrolyte s^{-1} (for run A2 of Packham), or $570 \text{ ml}^{-1}\text{s}^{-1}$ (Run A7). For the less restrictive run (A7), only at a triton energy of 14 KeV does the secondary neutron flux begin to fall below the threshold of detectability of neutrons (about 0.01s^{-1} and assuming an electrolyte content of 10 ml). We assume in these estimates that all other neutron channels (such as from $d+d \rightarrow {}^3\text{He}+n$) are totally suppressed.

If tritium is indeed created with the required energy of less than about 20 KeV, then one can expect an asymmetry of greater than 50 : 1 favoring the unwanted neutron branch in the dd fusion reaction. This arises as a simple consequence of the appearance in the transition matrix element of not only the entrance channel wave function (containing the usual tunneling amplitude) but also the exit channel wave function, which in the case of the charged reaction products $t+p$ being created at KeV energies will be suppressed due to the small amplitude of the Coulomb wave near to the nuclear channel radius in this exit channel. This effect is well known, such as in the Gamov description of α -radioactivity. Clearly, the ${}^3\text{He}+n$ channel does not suffer from this suppression, as there is no repulsive Coulomb force acting between the reaction products.

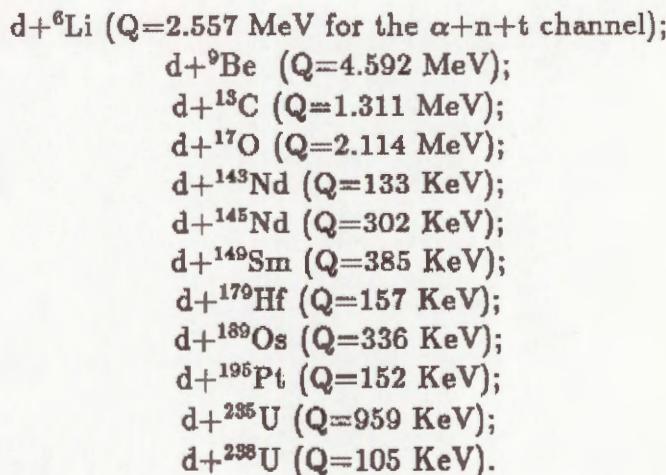
One may also wonder where the energy of fusion has gone, if tritons are formed at such small energies in the dd reaction. The only plausible answer is that such a reaction arises in association with the conversion of an electron, or a Bremsstrahlung photon. The phase space of the three final state particles favors equal momentum for each, but that implies that the momentum of the nuclei is only about 4 MeV/c. This corresponds to a very small energy for the triton, just a fraction of a KeV. While this is in accordance with the lack of secondary neutrons, such a tiny energy would lead to a strong branching into the neutron producing branch of the dd reaction because of the above-mentioned final state Coulomb effect. This observation is weakened to some degree by the possibility of (yet unexplained) long range neutron tunneling towards one of the deuterons, but one would expect such a process to produce energetic tritons.

Thus if indeed the tritium observed is due to cold fusion, we must find another nuclear reaction which will form tritons at just the "right" energy of about 20 KeV, higher energies being forbidden by absence of secondary neutrons, lower energies being made implausible by final state Coulomb suppression. We therefore turn now to consider all possible tritium producing, deuterium induced nuclear reactions, accessible in cold fusion. The number of direct exothermic

(positive Q-value) nuclear reactions involving a deuterium reacting with a stable or sufficiently long lived nucleus with nucleon number (A), resulting in tritium and a nucleus (A-1), is indeed severely limited; in fact, a survey of all such exothermic reactions using available nuclear and atomic mass tables [12] reveals just fifteen such possibilities, and two other candidates with small endothermic reaction energy (negative Q-values) which could possibly change sign after reevaluation of the experimental data. The two exothermic reactions:



are the only two known reactions with a *Q*-value of less than 100 KeV. The other 13 reactions, aside from the usual dd case, are:



None of these 13 more energetic reactions can proceed via a small Q-value to an excited intermediate state in the daughter nucleus. The nuclear mass data currently available further shows that $d + {}^{163}Dy$ has an adopted Q-value of $Q = -14 \pm 4$ KeV, having been listed at earlier times as $Q = 5$ KeV. Another case, $d + {}^{176}Lu$ has been listed with $Q = 66$ KeV in older mass tables, but today the adopted value is $Q = -36$ KeV.

Thus if tritium is produced without the presence of a much larger neutron signature, we must conclude that a (d,t) reaction on a heavy nucleus has taken place by some unknown mechanism with the best candidate being the reaction involving ${}^{201}Hg$ (13.6 natural relative abundance, resulting in a triton of 19 ± 6 KeV, just barely within the tolerable limit for neutron production by

secondary reactions. (Other such good candidates are ^{183}W (14.3% natural abundance) and perhaps also ^{168}Dy (24.9%).) In view of the obstacles to such a reaction and the preceding discussion we however conclude that the process of tritium production without accompanying neutron yield gives rise to even a greater scientific mystery as the purported anomalous heat effect.

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